

Professor: J.D. Wilson Time available: 120 mins Value: 35%

Please answer in the exam booklet. Symbols have their usual meteorological interpretation. Some data are given at the back.

Multi-choice**(22 x $\frac{1}{2}\%$ = 11%)**

- If the vector \vec{F} represents the convective flux density of a certain scalar property ϕ and \vec{u} the velocity field, then _____
 - $\vec{F} = \vec{u} \cdot \nabla \phi$
 - $\vec{F} \cdot \vec{u} = 0$
 - $\nabla \cdot (\vec{F} \vec{u}) = 0$
 - \vec{F} and \vec{u} are perpendicular
 - $\vec{F} = \vec{u} \phi$ ✓✓
- The units of a diffusivity (as appears, for example, in $\partial\phi/\partial t = K \partial^2\phi/\partial x^2$) are _____
 - m s^{-1}
 - m s^{-2}
 - $\text{m}^2 \text{s}^{-1}$ ✓✓
 - $\text{m}^2 \text{s}^{-2}$
 - kg m s^{-2}
- A continuous random variable q , defined on the range $-\infty \leq q \leq \infty$, belongs to a probability distribution whose probability density function is $f(q)$. If

$$\bar{q} \equiv \int_{-\infty}^{\infty} q f(q) dq = 0$$

then the variance of q is given by _____

- $\int_{-\infty}^{\infty} q^2 f(q) dq$ ✓✓
- $\int_{-\infty}^{\infty} f(q) dq \equiv 1$
- $\int_{-\infty}^{\infty} f(q) dq$
- $\int_{-\infty}^{\infty} dq$
- ∞

4. In the context of numerical solution of the advection equation

$$\frac{\partial \phi}{\partial t} + U \frac{\partial \phi}{\partial x} = 0$$

using finite differences in time (Δt) and space (Δx), a stability criterion on the Courant number arises. This criterion involves a ratio of velocities, one of which, namely _____, is _____ in origin

- (a) U ; numerical
 - (b) $\Delta x/\Delta t$; numerical ✓✓
 - (c) $\Delta x/\Delta t$; physical
 - (d) $\Delta t/\Delta x$; physical
 - (e) $\Delta \phi/\Delta x$; numerical
5. The Navier-Stokes equation (expressing conservation of momentum) and the conservation equation for a passive, non-reacting (“tracer”) species are given as data, as are conventions for representing the velocity vector. The advection term $u_j \partial c/\partial x_j$ in the tracer conservation equation may alternatively be written _____

- (a) $\vec{u} \cdot \nabla c$
 - (b) $\nabla (c\vec{u})$
 - (c) $u \frac{\partial c}{\partial x}$
 - (d) $u \frac{\partial c}{\partial x} + v \frac{\partial c}{\partial y} + w \frac{\partial c}{\partial z}$
 - (e) both a and d are valid ✓✓
6. Again referring (if necessary) to the Navier-Stokes equation and the tracer conservation equation (given as data, along with definitions of Reynolds, Rossby & Peclet numbers), the Lagrangian approach to treating dispersion of a tracer is valid in the limit of _____

- (a) infinitely small Reynolds number
- (b) infinitely large Reynolds number
- (c) infinitely small Peclet number
- (d) infinitely large Peclet number ✓✓
- (e) infinitely large Rossby number

7. According to the prescription covered in EAS471 (i.e. following Molinari, 1993), a “mesoscale” model of the atmosphere is _____, and has horizontal gridlengths in the range _____

- (a) non-hydrostatic; $10 \leq \Delta \leq 50$ km
- (b) hydrostatic; $10 \leq \Delta \leq 50$ km ✓✓
- (c) non-hydrostatic; $1 \leq \Delta \leq 10$ km
- (d) hydrostatic; $1 \leq \Delta \leq 10$ km
- (e) Eulerian; $\Delta \leq 10$ m

8. Let $\delta T(z)$, $\delta q(z)$ be adjustments to resolved temperature and specific humidity made in the layer $z_B \leq z \leq z_T$ of a particular grid column upon application of a cloud scheme(s) of the physics package of an NWP model. If these adjustments satisfy

$$\begin{aligned} -L \rho \delta q(z) &= \rho c_p \delta T(z), \quad (\delta T \geq 0), \\ q(z) + \delta q(z) &= q_*(T + \delta T, p), \end{aligned}$$

where $q_*(T)$ is the saturation specific humidity for temperature T , then the package that was activated corrects for _____

- (a) moist convection
 - (b) dry convection
 - (c) non-convective resolved-scale supersaturation (assumed to result in condensation) ✓✓
9. The two diagnostic conditions pertaining to the resolved state of the atmosphere and that cause the activation of the “Kuo scheme” for unresolved deep convection are _____ (where ∇_H is the horizontal grad operator, and \vec{v}_H the resolved horizontal velocity)
- (a) low level convergence $\nabla_H \cdot \vec{v}_H < 0$, and conditional instability of a deep layer ✓✓
 - (b) low level convergence $\nabla_H \cdot \vec{v}_H > 0$, and absolute instability of a deep layer
 - (c) low level convergence $\nabla_H \cdot \vec{v}_H < 0$, and absolute instability of a deep layer
 - (d) existence of CAPE (convectively available potential energy) exceeding 1000 J kg^{-1} , and supersaturation of a deep layer
10. The Kuo scheme, when activated, computes the quantity

$$M_t = - \int_0^\infty \nabla_H \cdot (\rho q \vec{v}_H) dz + E_0$$

where E_0 is the surface evaporation rate, q is resolved specific humidity, and other quantities are as above. This quantity represents _____

- (a) a supply of latent heat (and, ipso facto, water vapour) available to build convective clouds ✓✓
 - (b) the amount of precipitation produced by the scheme
 - (c) the energy used to raise the environmental temperature (i.e. resolved model state $\tilde{T}(z)$, pre-correction) to the temperature $T_c(z)$ of the diagnosed cloud
 - (d) local cloud fraction, i.e. fraction of sky occupied by deep convective cloud
11. The shortest wave that can be represented on a grid with spacing Δx has wavelength
- (a) $\pi/\Delta x$
 - (b) $2\pi/\Delta x$
 - (c) $\Delta x/2$
 - (d) Δx

(e) $2\Delta x$ ✓✓

12. The Canadian Meteorological Centre's Global Environmental Multiscale (GEM) model for Numerical Weather Prediction treats (resolved) advection terms by a "semi-Lagrangian" method. A fully Lagrangian approach to solving the momentum equations is not possible because _____

- (a) velocity (and momentum) are conserved properties
- (b) velocity (and momentum) *are not* conserved properties ✓✓
- (c) momentum is absorbed by the surface underlying the atmosphere
- (d) fluid elements are not absorbed by the surface
- (e) too many grid cells would be required

13. Suppose the distribution of a property $\phi = \phi(\mathbf{x}, t)$ in a fluid/gas is governed by

$$\frac{\partial \phi}{\partial t} = - \frac{\partial F_i}{\partial x_i} + Q(\mathbf{x}, t) ,$$

where $\mathbf{x} \equiv x_i$ signifies position (note: you might compare this " ϕ -eqn" with the species conservation equation given as data). The transport term in this equation is _____

- (a) $\partial \phi / \partial t$
- (b) Q
- (c) F_i
- (d) $-\partial F_i / \partial x_i$ ✓✓
- (e) $\int_{-\infty}^t \phi(\mathbf{x}, t) dt$

14. Suppose the distribution of a property $\phi = \phi(\mathbf{x}, t)$ in a fluid/gas is governed by

$$\frac{\partial \phi}{\partial t} = - \frac{\partial F_i}{\partial x_i} + Q(\mathbf{x}, t) ,$$

(the " ϕ -eqn") where $\mathbf{x} \equiv x_i$ signifies position. This is a _____ general statement than the species conservation equation given as data, in the sense that _____

- (a) more; the flux density F_i of the above ϕ -eqn is generic, embracing any or all transport mechanisms (i.e. potentially radiation, convection, diffusion/conduction), and a volumetric source term is included, allowing for in situ production or destruction of ϕ ✓✓
- (b) less; a volumetric source term is here included, allowing for in situ production or destruction of ϕ
- (c) less; no diffusion term appears in this equation for ϕ
- (d) less; no advection term appears in the ϕ equation

15. If ρ , T , c_p are the density, temperature and specific heat of air and \vec{u} is the velocity field, the convective flux density of sensible heat is _____

- (a) $\rho c_p \vec{u} \cdot T$
- (b) $\rho \vec{u} T$
- (c) $c_p \vec{u} T$
- (d) $\rho c_p \vec{u} T$ ✓✓
- (e) $\rho c_p \nabla \cdot (\vec{u} T)$

16. Suppose the profile of $f(x)$ along the x -axis is represented at discrete nodes, separated by interval Δx , and labelled J . If the derivative df/dx is represented as

$$\left(\frac{df}{dx}\right)_J = \frac{f_J - f_{J-1}}{\Delta x}$$

then the truncation error is of order

- (a) Δx ✓✓
- (b) Δx^2
- (c) $\sqrt{\Delta x}$
- (d) $(x_{J+1} + x_{J-1})/2$
- (e) $f_J - f_{J-1}$

17. Suppose $f(x, t)$ defined on $-1 \leq x \leq 1$, $0 \leq t$ is governed by the heat equation

$$\frac{\partial f}{\partial t} = K \frac{\partial^2 f}{\partial x^2}$$

with boundary conditions $f(-1, t) = f(1, t) = 0$, and that $f(x, 0) = \cos k\pi x/2$ where k is an arbitrary positive integer. Initially f has a maximum magnitude of 1, i.e. $|f|_{mx} = 1$. At later times the *true* state must satisfy

- (a) $|f|_{mx} = k$
- (b) $|f|_{mx} \geq 1$
- (c) $|f|_{mx} \leq 1$ ✓✓
- (d) $|f|_{mx} = \frac{K \Delta t}{\Delta x^2}$
- (e) $|f|_{mx} = \frac{k \Delta t}{\Delta x^2}$

18. Gridpoint computations for the influence of unresolved scales of motion in the ABL on the resolved absolute humidity $\bar{\rho}_v$ will involve the equation

$$\left(\frac{\partial \bar{\rho}_v}{\partial t}\right)_{\text{physics}} = - [\cdot] \overline{w' \rho'_v}$$

The missing operator “[.]” is _____

- (a) $U \partial/\partial x + V \partial/\partial y$ (U, V the resolved horizontal velocity components)
 - (b) $W \partial/\partial z$ (W the resolved vertical velocity)
 - (c) $\partial/\partial x$
 - (d) $\partial/\partial z$ ✓✓
 - (e) $K \partial/\partial z$ (K the eddy diffusivity)
19. The “curvature” of a scalar field $\phi(x, y, z)$ is given by (or measured by) application of the _____ operator
- (a) ∇
 - (b) $\nabla \cdot \nabla$ ✓✓
 - (c) $\hat{k} \cdot \nabla$
 - (d) $\hat{k} \times$
 - (e) $\hat{k} \times \nabla$
20. According to the Lax Equivalence Theorem, “If a difference equation is *consistent* with the differential equation it represents, then stability is the necessary and sufficient condition for convergence.” Here the technical meaning of “consistent” is that
- (a) truncation error must vanish in the limit of vanishing grid interval(s) ✓✓
 - (b) truncation error must vanish in the limit of infinite grid interval(s)
 - (c) the numerical solution ϕ^{num} equals the true (but generally unknown) solution ϕ to the differential equation in the limit of vanishing grid interval(s)
 - (d) the difference between the numerical solution ϕ^{num} to the difference equation and the (generally unknown) exact solution ϕ^* to the difference equation vanishes in the limit of vanishing grid interval(s)
21. According to the Lax Equivalence Theorem, “If a difference equation is consistent with the differential equation it represents, then stability is the necessary and sufficient condition for *convergence*.” Here the technical meaning of “convergence” is that
- (a) truncation error must vanish in the limit of vanishing grid interval(s)
 - (b) truncation error must vanish in the limit of infinite grid interval(s)
 - (c) the numerical solution ϕ^{num} equals the true (but generally unknown) solution ϕ to the differential equation in the limit of vanishing grid interval(s) ✓✓
 - (d) the difference between the numerical solution ϕ^{num} to the difference equation and the (generally unknown) exact solution ϕ^* to the difference equation vanishes in the limit of vanishing grid interval(s)

22. The Random Displacement Model (RDM) for the vertical motion of a fluid element (or “particle”) is

$$dZ = \frac{\partial K}{\partial z} dt + \sqrt{2K} dt r ,$$

where dZ denotes the increment in height during the time step dt , K is the eddy diffusivity, and r is chosen from a Normal distribution with zero mean and unit variance. The RDM

- (a) is an Eulerian method
- (b) requires the imposition of a spatial grid in order to allow computation of trajectories
- (c) is a valid description of dispersion even in the near field of a source (i.e. for travel times t much smaller than the Lagrangian time scale for the vertical velocity)
- (d) is a grid free (Lagrangian) treatment of turbulent convection that is *equivalent* to the “diffusion” model ✓✓
- (e) is also known as the “generalized Langevin equation” or the “first-order Lagrangian stochastic model”

Short answer:

3 x 8% = 24%

Answer any **three** questions from this section.

1. Perform a dimensional analysis to find the form of the law for the drag force F on a sphere of density ρ and radius R that is falling at velocity V through still air whose density and kinematic viscosity are respectively ρ_a, ν (the units of kinematic viscosity are $\text{m}^2 \text{s}^{-1}$).
2. Determine the 4×4 tridiagonal coefficient matrix \mathbf{M} and the right hand side \mathbf{B} in a matrix expression of form $\mathbf{M} \Theta = \mathbf{B}$ for the numerical solution of

$$\frac{\partial^2 \theta}{\partial z^2} = 2\theta$$

on $0 \leq z \leq 1$, subject to $\theta(0) = 0, \theta(1) = 1$. Set up your solution with four, equi-spaced gridpoints indexed $J = (1, 2, 3, 4)$ positioned at $z_J = (0, 1/3, 2/3, 1)$. At internal gridpoints (J=2,3) the curvature is to be represented as

$$\frac{\partial^2 \theta}{\partial z^2} = \frac{\theta_{J+1} + \theta_{J-1} - 2\theta_J}{\Delta z^2} .$$

Note: you are not being asked to invert \mathbf{M} , nor to obtain the (numerical) solution vector $\Theta = (\theta_1, \theta_2, \theta_3, \theta_4)$.

3. Briefly summarize the Canadian Meteorological Centre’s Global Environmental Multiscale (GEM) model of the atmosphere, as used for short range (48 hour) Numerical Weather Prediction. Your response should cover salient points in regard both to model dynamics and model physics (grid point computations).

4. In the context of grid point computations, which treat the atmospheric boundary layer as if it were horizontally-homogeneous, the conservation equation for the kinetic energy $k \equiv \frac{1}{2}(\overline{u'^2} + \overline{v'^2} + \overline{w'^2})$ residing in the unresolved scales of motion can be approximated as:

$$\frac{\partial k}{\partial t} = K \left[\left(\frac{\partial \bar{u}}{\partial z} \right)^2 + \left(\frac{\partial \bar{v}}{\partial z} \right)^2 \right] - \frac{g}{\theta_0} \frac{K}{P_r} \frac{\partial \bar{\theta}}{\partial z} - \epsilon + \frac{\partial}{\partial z} \left(K_k \frac{\partial k}{\partial z} \right),$$

where θ_0 is the mean potential temperature of the layer, (\bar{u}, \bar{v}) and $\bar{\theta}$ are the resolved components of the horizontal wind and the potential temperature, K is the eddy viscosity, $K_h = K/P_r$ is the eddy diffusivity for heat (P_r is the turbulent Prandtl number), and K_k is the effective eddy diffusivity for the vertical diffusion of k .

Classify each term in this equation, and identify its conventional name — e.g. what term(s) represent “shear production”? Making reference to the flux Richardson number

$$R_i^f = \frac{g}{\theta_0} \frac{K/P_r}{K} \frac{\partial \bar{\theta}/\partial z}{(\partial \bar{u}/\partial z)^2 + (\partial \bar{v}/\partial z)^2},$$

explain the influence of atmospheric stratification on the energy of the unresolved motion (you may assume $P_r = 1$).

5. Assuming a hydrostatic atmosphere, the (vertical) vorticity equation is

$$\begin{aligned} \frac{\partial \zeta}{\partial t} + \vec{v}_H \cdot \nabla_H (\zeta + f) + \omega \frac{\partial \zeta}{\partial p} \\ = - (\zeta + f) \nabla_H \cdot \vec{v}_H + \hat{k} \cdot \left(\frac{\partial \vec{v}_H}{\partial p} \times \nabla_H \omega \right) \end{aligned} \quad (1)$$

where ζ is the relative vorticity, \vec{v}_H is the horizontal wind vector and ω [Pa s⁻¹] the vertical velocity. Explain the further limitations/assumptions/substitutions and/or simplifications that lead to the quasi-geostrophic vorticity equation

$$\frac{d_g \zeta}{dt} \equiv \frac{\partial \zeta}{\partial t} + U_g \frac{\partial \zeta}{\partial x} + V_g \frac{\partial \zeta}{\partial y} + V_g \beta = f_0 \frac{\partial \omega}{\partial p}$$

where \vec{V}_g is the Geostrophic wind vector defined at a reference latitude ϕ_0 , f_0 is the Coriolis parameter at that latitude, and $\beta = (\partial f / \partial y)_{\phi_0}$ its northerly gradient.

Data

- $(\hat{i}, \hat{j}, \hat{k})$ are unit vectors along the Cartesian coordinate directions (x, y, z) , and (u, v, w) are the corresponding Cartesian velocity components. Alternative notations for the velocity vector are \mathbf{u} , \vec{u} and u_i , where the dummy subscript occurring in the last of these (and which could as easily have been written j or k or indeed any other symbol) takes on values of 1,2 or 3 corresponding to the three spatial axes.

•

$$\frac{d\mathbf{u}}{dt} \equiv \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u} = \frac{-1}{\rho} \nabla p + \mathbf{g} - 2\boldsymbol{\Omega} \times \mathbf{u} + \nu \nabla^2 \mathbf{u}$$

The Navier-Stokes equation, expressing conservation of momentum of a Newtonian fluid: ρ the fluid density, p the pressure, \mathbf{g} the gravitational acceleration vector, $\boldsymbol{\Omega}$ the angular velocity of the coordinate frame (occurring in the Coriolis term), ν the molecular kinematic viscosity. If V, L are velocity and length scales for the motion, a Reynolds number may be formed as $R_e = VL/\nu$ and molecular friction can be neglected in the limit $R_e \rightarrow \infty$. Similarly the Rossby number is $R_o = V/(fL)$ where $f = 2|\boldsymbol{\Omega}| \sin \phi$ is the Coriolis parameter.

•

$$\frac{dc}{dt} \equiv \frac{\partial c}{\partial t} + u_j \frac{\partial c}{\partial x_j} = \kappa \frac{\partial^2 c}{\partial x_j \partial x_j}$$

Expresses conservation of the mass of a passive, non-reactive gas in the atmosphere. The velocity vector $\mathbf{u} = u_i$, and summation applies in any term with a repeated subscript, e.g. $u_j \partial c / \partial x_j \equiv \vec{u} \cdot \nabla c$. The molecular diffusivity of the gas in air is κ . If V, L are velocity and length scales for the motion, a Peclet number may be formed as $P = VL/\kappa$

- $D_p + \frac{\partial \omega}{\partial p} = 0$

Continuity equation in the x, y, p (“isobaric”) coordinate system, where $D_p \equiv \nabla_H \cdot \vec{v}_H$ is the divergence of the two velocity components lying in the constant pressure surface (“horizontal divergence”)

Schematic answers

1. There are $n = 6$ variables involved¹ ($F, \rho, R, V, \rho_a, \nu$) and there are $m = 3$ fundamental dimensions (mass, length, time). Thus we seek a law of form

$$\pi_1 = f(\pi_2, \pi_3)$$

where π_1 will be the ratio of the force F to some suitable scale for force, such as $\rho_a V^2 R^2$ (force = pressure \times area, and density \times velocity² has units of pressure). By inspection, one of the ratios in the unknown function can be chosen as $\pi_1 = \rho/\rho_a$. Intuition demands that the other ratio should involve the viscosity (for if there were no viscosity, the potential exists for the sphere to fall without drag). Thus we may choose $\pi_2 = RV/\nu$, which gives the latter the form of a Reynolds number. With these choices, the law we seek is

$$\frac{F}{\rho_a V^2 R^2} = f\left(\frac{\rho}{\rho_a}, \frac{RV}{\nu}\right).$$

Other correct results exist. For example one student used the method of indices to find a scale for force, with the result that π_1 was chosen as

$$\pi_1 = \frac{F}{\rho^2 R V \rho_a^{-1} \nu}.$$

This is perfectly correct, albeit more complex than it needs to be. Other correct scalings for π_1 are the choices:

$$\begin{aligned} \pi_1 &= \frac{F}{\rho_a^2 R V \rho^{-1} \nu}, \\ \pi_1 &= \frac{F}{\rho R V \nu}, \\ \pi_1 &= \frac{F}{\rho_a R V \nu}. \end{aligned}$$

Table 1: Scoring criteria to maximum of 8.

diagram	+1
recognizing $n = 6$	+1
recognizing $m = 3$	+1
seeking $n - m$ ratios π (irrespectively of whether n, m were assigned correctly)	+2
for each correct π	+2

¹Note that the gravitational acceleration g was not mentioned, and it was not stated that the sphere is falling at a steady rate, i.e. at equilibrium. Occurrence of the word “falling” (rather than “moving”) betrays a subconscious notion of existence of vertical motion under gravity, and so it would have been acceptable if students were to include g and conclude $n = 7$.

2. The set-up for numerical solution goes as follows². We label the four values at the grid-points $\theta_1 \dots \theta_4$. The gridlength $\Delta z = 1/3$. We wish to encode into a matrix the four equations:

$$\begin{aligned}\theta_1 &= 0 \\ \theta_1 - 2\theta_2 + \theta_3 &= 2\theta_2\Delta z^2 \\ \theta_2 - 2\theta_3 + \theta_4 &= 2\theta_3\Delta z^2 \\ \theta_4 &= 1\end{aligned}$$

(note the simplicity of the boundary conditions). Thus the matrix formulation of the problem is

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 1 & -2 - 2\Delta z^2 & 1 & 0 \\ 0 & 1 & -2 - 2\Delta z^2 & 1 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \theta_1 \\ \theta_2 \\ \theta_3 \\ \theta_4 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix},$$

where $-2 - 2\Delta z^2$ evaluates to $-2 - 2/9 = -20/9$. Note that one could multiply all coefficients by -1 , in any of rows (1,2,3), without changing the answer.

Table 2: Scoring criteria to maximum of 8.

relevant diagram	+1
giving a tridiagonal <i>form</i> for \mathbf{M}	+1
correct boundary conditions	+1 \times 2
correct interior equations	+2 \times 2

²As an aside — students were not asked to do this — the analytical solution $\theta(z)$ is easy to obtain. By substituting a trial solution $\theta = e^{mz}$ into the differential equation. It follows immediately that $m^2 = 2m$ and so $m = \pm\sqrt{2}$ and the general solution is

$$\theta = ae^{\sqrt{2}z} + be^{-\sqrt{2}z}.$$

By using the boundary conditions the constants a, b are easily determined, and the solution is

$$\theta(z) = 0.258 \left[e^{\sqrt{2}z} - e^{-\sqrt{2}z} \right].$$

3. Valid points relating to GEM in its configuration for the Regional runs (to maximum of 8%).

Table 3: Valid points.

Dynamics	Physics (parameterizations)
Hydrostatic (1) Primitive eqns for u, v (1) Global domain (1) Timestep 7.5 min (1) Horiz. gridlength 15 km (1) Terrain-following vertical coordinate (1) Highest level 10 hPa ($\frac{1}{2}$) Nearly 60 levels ($\frac{1}{2}$) Nominally 8 levels in ABL ($\frac{1}{2}$) Run 4× daily ($\frac{1}{2}$) Semi-Lagrangian treatment of advection ($\frac{1}{2}$) Non-uniform Δ_x, Δ_z outside N. Amer. ($\frac{1}{2}$)	Gdpnt compt'ns assume (local) horiz. homogeneity (1) Atmos. bndry layer (ABL) treated by eddy diffusion (1) Deep convection/precip ($\frac{1}{2}$) Shallow convection/precip ($\frac{1}{2}$) Stratiform cloud/precip ($\frac{1}{2}$) Solar radiation ($\frac{1}{2}$) Longwave radiation ($\frac{1}{2}$) Gravity waves ($\frac{1}{2}$) Sfc types incl. canopy, soil, water/snow/ice ($\frac{1}{2}$)

4. Conservation equation for the kinetic energy k of the unresolved motion field, and interpretation of terms so as to explain the influence of stratification on k .

Background: If the atmosphere is “stratified,” then by definition there exists a vertical gradient in potential temperature. The usual corollary of that gradient is a vertical flow of heat, carried by the unresolved motion — except in the case of extremely stable stratification, when the unresolved motion ceases to occur, and $k = 0$.

- Terms are classified as follows:

$$\frac{\partial k}{\partial t} = K \left[\left(\frac{\partial \bar{u}}{\partial z} \right)^2 + \left(\frac{\partial \bar{v}}{\partial z} \right)^2 \right] - \frac{g}{\theta_0} \frac{K}{P_r} \frac{\partial \bar{\theta}}{\partial z} - \epsilon + \frac{\partial}{\partial z} \left(K_k \frac{\partial k}{\partial z} \right)$$

storage
source
source
source
transport

- The shear production terms are

$$P_S = K \left[\left(\frac{\partial \bar{u}}{\partial z} \right)^2 + \left(\frac{\partial \bar{v}}{\partial z} \right)^2 \right]$$

- Buoyant production is

$$P_B = - \frac{g}{\theta_0} \frac{K}{P_r} \frac{\partial \bar{\theta}}{\partial z},$$

and vanishes if the potential temperature is height-independent. Expressed more fundamentally, this term would explicitly show the vertical heat flux, viz. be $P_B \equiv (g/\theta_0) \overline{w'\theta'}$, but the introduction of K-theory to parameterize the heat flux results in the form shown

- The turbulent kinetic energy dissipation term is ϵ . This term is non-negative, and is always a loss term, tending to damp down the unresolved scales of motion
- The “transport term” is

$$P_T = \frac{\partial}{\partial z} \left(K_k \frac{\partial k}{\partial z} \right)$$

- The flux Richardson number is $R_i^f = P_B/P_S$ and we notice that if the turbulent Prandtl number $P_r = 1$ then R_i^f simplifies to what is called the “gradient” Richardson number

$$R_i^g = \frac{g}{\theta_0} \frac{\partial \bar{\theta} / \partial z}{(\partial \bar{u} / \partial z)^2 + (\partial \bar{v} / \partial z)^2}.$$

We note that the energy of the unresolved motion is characterized by (measured by) the property k . Physically, the role of stratification is as follows:

- If the layer is neutral, the Richardson numbers are zero. Buoyancy forces do not occur in the flow, and the unresolved motion is “powered” by shear production alone

- In an unstable layer $\partial\bar{\theta}/\partial z < 0$ and so $P_B > 0$, and buoyancy acts to enhance the kinetic energy of the unresolved motion field. The Richardson numbers are negative.
- In a stable layer $\partial\bar{\theta}/\partial z > 0$ and so $P_B < 0$, and buoyancy acts to reduce the kinetic energy of the unresolved motion field. The Richardson numbers are positive.

Table 4: Scoring criteria to maximum of 8.

recognizing terms by type (storage/transport/source)	+2
recognizing terms by conventional name (“shear produc’n” etc.)	+2
recognizing the flux Richardson number as ratio of buoyant to shear production	+2
associating “stratification” with $\partial\bar{\theta}/\partial z \neq 0$	+1
associating “neutral stratification” with $\partial\bar{\theta}/\partial z = 0$	+1
associating “neutral stratification” with $\overline{w'\theta'} = 0$	+1
giving the flux-gradient relationship $\overline{w'\theta'} = -K_h \partial\bar{\theta}/\partial z$	+1
associating the potential temperature gradient $\partial\bar{\theta}/\partial z$ with a vertical heat flux	+1
associating the sign of the Richardson number with stability/neutrality/instability	+2
associating non-zero R_i with suppression/enhancement of k	+2

5. Assumptions leading to the quasi-geostrophic vorticity equation. This question was given on the 2007 exam — see answer given there. This year no student chose to answer this question.